

# Passive Wireless Vibration Sensors Based on the Amplitude Modulation of Harmonic Backscattering

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**Abstract**—This article presents a wireless passive vibration sensor based on harmonic backscattering. The adopted vibration transducer is a commercial piezoelectric cantilever. The cantilever is connected to the gate of an NMOS transistor, which controls the self-bias point of the Schottky diode responsible for the second-harmonic generation. This way, vibrations induce the amplitude modulation of the backscattered second harmonic. A technique is presented to retrieve the vibration information (both acceleration and frequency) from the acquired spectrum. The vibration frequency is computed as the average frequency difference between the carrier and the first sidebands, while the vibration acceleration is encoded in the duty cycle of the amplitude modulation of the backscattered harmonic. In this contribution, we demonstrate for the first time that this information can be retrieved from the magnitude ratio of the first and second sidebands of the acquired spectrum regardless of the tag-to-reader distance. A complete wireless prototype, working at  $f_0 = 1.5$  GHz, is manufactured and tested in a laboratory environment for tag-to-reader distances up to 1.3 m. Despite some sensor limitations, i.e., the nonlinear relationship between the duty cycle and the vibration acceleration and the dependence of the sensor sensitivity on the vibration frequency, excellent results are obtained for vibration frequencies around the natural oscillation frequency of the cantilever (i.e., 130 Hz), and for variable accelerations and distances, which testifies the high sensitivity and the robustness of the proposed solution.

**Index Terms**—Backscatter radio, harmonic backscattering, harmonic transponder, Internet of Things, piezoelectric cantilever, vibration measurement, vibration sensor, wireless sensor.

## I. INTRODUCTION

INTERNET of Things (IoT) solutions are more and more used in industrial applications to perform condition monitoring and predictive maintenance of industrial equipment, with the aim to optimize resource usage, reduce waste, and

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improve safety [1]. Wireless sensors, in particular, are the backbone of IoT systems, as they conveniently provide the required real-time data to decision-making algorithms and to dynamic digital models [2], [3].

Among wireless sensors for industrial applications, vibration sensors play a pivotal role. They are used to perform structural health monitoring (SHM), which is critical to predict risks by checking the wear status of equipment and infrastructures and the stress that mechanical parts are subject to [4]. Vibration sensors are based on a variety of transducers, ranging from piezoelectric and piezoresistive materials [5], [6], to micro-electro mechanical (MEMS) accelerometers [7], [8], strain gauges [9], fiber optics [10], and variable capacitors [11].

The sensor node usually includes the vibration transducer and an active signal conditioning circuit, meaning that an external energy source (usually a battery) is required to power the active circuitry [4]. In addition, energy is needed for wireless communication. However, battery-operated wireless vibration sensors are inconvenient. The weight of the battery can alter the dynamic behavior of the structure to be monitored. In addition, batteries must be replaced periodically, which is an issue for sensors not directly accessible. This is usually the case of vibration sensors, as they can be installed on large tall structures, and inside industrial equipment. In the latter case, periodical battery replacements may cause direct interaction with the monitored equipment and additional costs due to the involvement of dedicated personnel.

For these reasons, research is focusing on developing low-power or even passive wireless vibration sensors [12]. In this context, wireless transponders based on backscatter radios are currently being developed [13], [14], [15]. Backscatter radios are interrogated by wireless signals generated by a reader and they modulate the reflected signals based on the information to be conveyed. This way, they can communicate wirelessly without generating any radio signal on board, which drastically reduces the power consumption for the wireless node [16]. Several attempts are reported in the literature to make backscatter radios energy autonomous, and to extend their read range [17], [18], while making them monitor physical parameters of interest [19]. Additional problems of backscatter radios are self-jamming and self-interference [20], [21], [22], and a few studies are reported, aimed at shifting

the spectrum of the backscattered signal to different frequency bands with respect to the transmitted carrier [23], [24].

A few vibration sensing solutions have been explored based on various types of backscatter radios [25], [26], [27], [28], [29]. In [25], a wireless vibration sensor is proposed based on a strain gauge and a voltage-controlled oscillator (VCO). The strain gauge determines the output voltage of an operational amplifier used to control the oscillation frequency of the VCO (the oscillation frequency is proportional to the strain signal), which, in turn, controls the impedance switching circuit used to modulate the backscattered signal. Therefore, the vibration frequency is encoded as the frequency of the amplitude-modulated backscattered signal. Analogously, the wireless sensor in [26] is based on a microstrip patch antenna and an amplitude modulator powered by a light-power energy harvester, where the resonant frequency of the antenna is a function of the tensile strain applied. However, both the circuit solutions need an additional energy source to operate, which is provided by a solar energy harvester, and therefore, they might not be suitable for environments with dim light conditions, although solar cells optimized for indoor applications are currently being studied [30].

In [27], a surface acoustic wave (SAW) passive wireless vibration sensor is proposed, based on an Y-cut quartz cantilever beam with a proof mass at the free end. Vibrations flex the cantilever, inducing relative changes in the acoustic propagation through the SAW resonators deposited on the surface of the beam and causing a frequency variation in the generated output electromagnetic signal. However, SAW tags are usually more expensive than transponders based on the traditional printed circuit boards (PCBs) and commercial off-the-shelf components and can have limited read range [31].

The wireless passive vibration sensor illustrated in [28], instead, is an LC resonant circuit based on a constant square inductor and a variable parallel-plate capacitor with a movable plate. The LC circuit is magnetically coupled to the reader coil. Under vibration, the capacitance changes due to the motion of the movable plate, which changes the reflection coefficient at the input of the reader coil. In this solution, though, the allowed tag-to-reader distance is extremely limited (in the order of 2 mm).

Among backscatter radios, harmonic transponders are attracting growing interest. Harmonic transponders convert the interrogating signal into its second harmonic, which is then backscattered. Since the backscattered signal lies in a different band with respect to the interrogating signal, the sensitivity of the receiver is not affected by eventual leakage of  $f_0$  signal from the transmitter, or by signal reflections from the environment, like in traditional backscatter radios. Therefore, high receiver sensitivities can be achieved [32]. The sensitivity of harmonic receivers can be impaired by the spurious second harmonic of the transmitter; however, this contribution can be minimized by adding a low-pass filter between the signal generator and the transmitting antenna. Harmonic transponders able to work with radio frequency (RF) input powers as low as  $-30$  dBm have been designed and demonstrated [33]. The higher sensitivity of the receiver and of the tag compensate for the power loss caused by the frequency conversion (variable

contribution, depending on the RF input power, the diode nonlinearity, and the circuit status [34]) and the higher path loss experienced by the backscattered second harmonic with respect to the fundamental tone (6-dB additional loss), leading to wireless sensing systems with potential longer read ranges than the traditional backscatter radios, especially in highly cluttered environments. In addition, their inherent immunity to clutter makes harmonic transponders good candidates for high-sensitivity sensors, and currently research has focused on adding sensing capabilities to harmonic transponders [35], [36], [37], [38], [39].

In [29], two circuit implementations for vibration sensing are presented based on harmonic backscattering. In the first circuit, a piezoelectric cantilever is used to change the bias point of a GaAs diode, leading to an amplitude modulation of the output signal. However, the impedance seen by the cantilever is lower than the optimal one, which leads to poor sensitivity. In the second circuit type, a mechanical switch is used to open and close the diode dc current path at the doubler input. Such a circuit can be used to detect low-frequency vibrations; however, the vibration acceleration cannot be measured.

In the present study, a wireless passive vibration sensor based on harmonic backscattering is presented, where a piezoelectric cantilever is used to switch an NMOS transistor between OFF and triode region. Both vibration acceleration and frequency can be recovered from the backscattered signal regardless of the tag-to-reader distance, using a mix of the two approaches described in [29]. This article is organized as follows. The theory of operation is shown in Section (II). Section III is dedicated to discuss the experimental results, and finally the conclusions are drawn.

## II. THEORY OF OPERATION

This article describes a wireless passive vibration sensor suitable for condition monitoring. The transponder relies on harmonic backscattering: a reader interrogates the tag using a sinusoidal signal with frequency  $f_0$ , while the tag backscatters the signal at the second harmonic [33]. The sensor induces the amplitude modulation of the harmonic backscattered signal, without requiring any additional energy source. Thanks to the frequency separation between the transmitted and received signals, harmonic transponders are robust to clutter, which makes them preferable in industrial environments [40].

The key component of the proposed transponder is a passive frequency doubler, which is shown in Fig. 1. The frequency doubler is a two-port network which aims at converting the RF input signal at the fundamental frequency  $f_0$  into its second harmonic at  $2f_0$ . It is composed of a series-connected low-barrier Schottky diode (nonlinearity), two quarter-wave stubs at  $f_0$  (harmonic filters) connected in shunt at the two sides of the diode, and the input and output matching networks (the circuit is designed for source and load impedances of  $50 \Omega$ ).

In the state-of-the-art frequency doubler in [41], the output matching network is composed of an LC network, where the inductor (see  $L_m$  in Fig. 1) and the capacitor (see  $C_m$  in Fig. 1) also provide a return path for the self-generated dc component. Indeed, when a sinusoidal signal is applied to a Schottky diode,

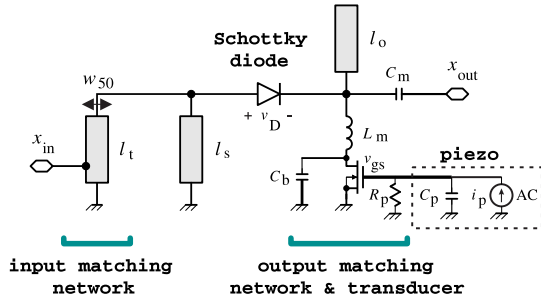


Fig. 1. Schematic of the proposed frequency doubler. The main parameters are:  $w_{50} = 1.62$  mm,  $l_t = 25.8$  mm,  $l_s = 29.7$  mm,  $l_s = l_o = 29.7$  mm,  $C_m = 0.4$  pF,  $L_m = 4.6$  nH,  $C_b = 50$  pF, and  $C_p = 22$  nF.

a dc component is generated along with the higher harmonic components [33], [34]. The series capacitor  $C_m$  prevents the dc component from entering the output load. On the other hand, the short-circuited harmonic stub, connected to the diode anode, and the inductor  $L_m$ , connected to ground, form a closed path where the self-generated dc current can circulate, while the dc voltage across the diode is forced to zero, i.e., the diode is zero-biased.

In the present modified version of the frequency doubler, a low-frequency NMOS transistor switch is connected between the inductance  $L_m$  and ground. The gate of the transistor is used to modulate the amplitude of the generated second-harmonic signal (hence called “modulation port”), by controlling the only dc component generated by the diode [34]. A bypass capacitor  $C_b$  is connected in parallel between the drain terminal of the transistor and ground, so that the transistor state only affects the self-generated dc component ( $L_m$  is RF-connected to ground). This way, the switch parasitics do not affect the behavior of the frequency doubler.

When the gate–source voltage  $V_{gs}$  is above the threshold voltage of the transistor  $V_{th}$ , the transistor is in the triode region and the channel is formed. Therefore, the inductance is dc-connected to ground and the Schottky diode is zero-biased (ON state); in this case, the maximum RF power is generated and transmitted to the output port of the frequency doubler.

When the gate–source voltage is below threshold, the channel is not formed instead,  $L_m$  is not dc-connected to ground, and the dc current is zero. In this case, it can be demonstrated that a positive voltage is developed at the diode cathode, and that the diode is reverse-biased. In particular, the dc voltage across the diode  $v_{D,dc}$  can be expressed as a function of the diode voltage at the fundamental frequency  $v_{D,f0}$  as follows [34]:

$$v_{D,dc} = -nV_T \ln \left[ I_0 \left( \frac{v_{D,f0}}{nV_T} \right) \right] \quad (1)$$

where  $I_0$  is the modified Bessel function of the first kind of the zeroth order,  $n$  is the diode ideality factor, and  $V_T$  is the thermal voltage. Since  $I_0(x)$  is a positive monotonically increasing function of  $x = v_{D,f0}/(nV_T)$ ,  $v_{D,dc}$  is negative, meaning that the diode is reverse-biased. In this condition, the power of the generated second harmonic is significantly reduced (OFF state). It is worth noting that in the proposed circuit the diode is never forward-biased, since no external bias

is applied to the diode. Therefore, the frequency doubler has the minimum conversion loss when the diode is zero-biased.

A piezoelectric cantilever is connected to the NMOS transistor to implement a passive wireless vibration sensor, as shown in Fig. 1. The piezoelectric cantilever can be modeled as a current source  $i_p$  in parallel with a capacitor  $C_p$  [5], [42]. Under vibration, the piezoelectric cantilever develops a time-varying voltage  $v_{gs}$  between the gate and the source of the MOSFET. For simplicity, sinusoidal vibrations are considered in this analysis, which leads to a sinusoidal voltage signal generated by the piezoelectric cantilever. The frequency of the generated signal depends on the vibration frequency  $f_v$ , and, for a given vibration frequency, its amplitude  $V_{gs}$  depends on the vibration acceleration  $a_v$  (the higher the vibration acceleration, the higher the voltage generated by the piezoelectric cantilever). Therefore, the voltage signal at the gate of the transistor can be expressed as follows:

$$v_{gs}(t) = V_{gs} \cos(2\pi f_v t + \phi_0) \quad (2)$$

where  $\phi_0$  is the initial phase of the signal. Thanks to the high impedance at the gate port of the transistor ( $C_{gs} \approx 50$  pF), a very small power is needed to switch ON and OFF the NMOS transistor switch [34], which can be easily provided by the piezoelectric cantilever [43]. A  $1$  M $\Omega$  resistor ( $R_p$  in Fig. 1) is connected in parallel with the cantilever to have a path to periodically discharge the capacitor  $C_{gs}$  based on the voltage applied by the cantilever.

When the voltage signal applied to the gate of the transistor is above a given threshold voltage, the switch is closed and the frequency doubler is in the ON state, while when the voltage is below threshold (including negative voltages) the transistor switch is open and the frequency doubler is in the OFF state. The fraction of the vibration period in which the frequency doubler is ON or OFF depends on the threshold voltage  $V_{th}$  and ON the amplitude of the voltage signal  $V_{gs}$ , as shown by Fig. 2(a). Let us define the duty cycle  $D$  of the amplitude-modulated second harmonic as follows:

$$D = \frac{T_{ON}}{T_v} \quad (3)$$

where  $T_{ON}$  is the time interval in which the transistor switch is ON, and  $T_v = 1/f_v$  is the vibration period.

Sinusoidal signals with larger amplitudes, corresponding to larger accelerations, lead to larger duty cycles. Indeed,  $D$ ,  $V_{gs}$ , and  $V_{th}$  are related as follows:

$$D = \frac{1}{\pi} \arccos \left( \frac{V_{th}}{V_{gs}} \right). \quad (4)$$

On a first approximation, the output signal  $x_{out}$  can be modeled as the second harmonic generated by the diode multiplied by a square wave with period  $T_v$  and duty cycle  $D$

$$x_{out} = \cos(2\omega_0 t) \left[ A_{avg} + (A_{ON} - A_{OFF}) \sum_{n=1}^{\infty} a_n \cos(n\omega_v t) \right] \quad (5)$$

where

$$A_{avg} = A_{ON} D + A_{OFF} (1 - D). \quad (6)$$

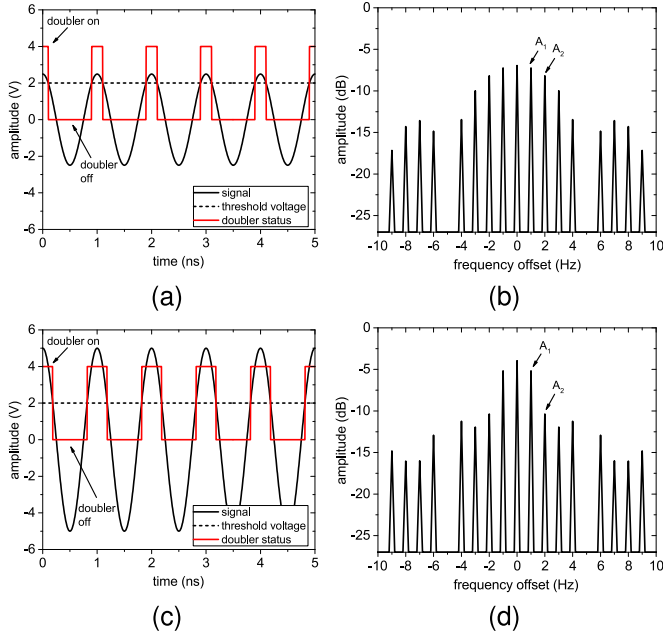


Fig. 2. Dependence of the duty cycle of the amplitude-modulated second harmonic on the amplitude of the signal generated by the piezoelectric transducer: (a) and (b) 20% duty cycle and (c) and (d) 40% duty cycle. (a) and (c) Time-domain signals generated by the piezoelectric cantilever and (b) and (d) output spectra. For simplicity, (b) and (d) are generated assuming  $A_{ON} = 1$  V,  $A_{OFF} = 0$  V, and the  $x$ -axis is normalized with respect to the carrier frequency.

$A_{ON}$  is the amplitude of the RF output signal when the switch is closed (the so-called “ON” state),  $A_{OFF}$  is the amplitude of the RF output signal when the switch is open (the so-called “OFF” state),  $\omega_0 = 2\pi f_0$ ,  $\omega_v = 2\pi f_v$ , and

$$a_n = \frac{2}{n\pi} \sin(n\pi D) \quad (7)$$

are the Fourier coefficients of the square wave.

The Fourier transform of  $x_{out}$  corresponds to the convolution of the Fourier transform of the sinusoidal RF signal and the square wave, as follows:

$$X_{out}(f) = A_0\delta(f - 2f_0) + \sum_{n=1}^{\infty} A_n\delta(f - 2f_0 \pm nf_v) \quad (8)$$

where

$$A_0 = \frac{A_{avg}}{2} \quad (9)$$

and

$$A_n = \frac{(A_{ON} - A_{OFF})a_n}{4}. \quad (10)$$

Generally, we need both magnitude and phase of the Fourier transform of  $x_{out}$  to retrieve the time-domain signal, since the spectral lines associated with the generated sidebands are multiplied by coefficients  $a_n$  that can be positive or negative based on the sign of the term  $\sin(n\pi D)$ .

In the present contribution, a method is proposed to retrieve the RF output signal  $x_{out}$  from the only spectrum acquired by the receiver. This method leverages the fact that the duty cycle of the square wave is always below 50% (the sinusoidal voltage generated by the piezoelectric cantilever is negative

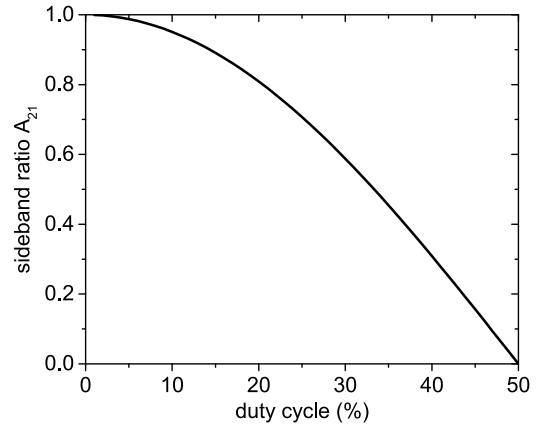


Fig. 3. Ratio of the magnitude of the second sideband (component at  $f_0 \pm 2f_v$ ) to the first sideband (component at  $f_0 \pm f_v$ )  $A_{21}$  versus the duty cycle of the square wave  $D$  for practical  $0 < D < 0.5$ .

for half of its period, which implies it is below the threshold voltage of the transistor). Therefore, the coefficients associated with the first and second sidebands ( $A_1$  and  $A_2$ ) are necessarily positive, and their ratio  $A_{21}$  is equal to

$$A_{21} = \frac{A_2}{A_1} = \frac{a_2}{a_1} = \frac{1 \sin(2\pi D)}{2 \sin(\pi D)}. \quad (11)$$

For  $D < 0.5$ ,  $D \neq 0$ , this is a positive, monotonically decreasing function of  $D$ , as shown in Fig. 3. Therefore, we can associate a unique  $D$  value with each  $A_{21}$  value, and  $A_{21}$  can be used to measure the vibration acceleration. Fig. 2(b) and (d) shows the spectrum associated with a duty cycle of 20% and 40%, respectively, clearly indicating the components  $A_1$  and  $A_2$ . It can be noted that  $A_{21}$  does not depend on  $A_{ON}$  and  $A_{OFF}$ ; therefore, this parameter is independent of the RF input power, which makes the proposed transducer and acquisition method generically suitable for wireless transponders.

Once  $D$  is known, the sign of all the sidebands can be calculated using (7), and the time-domain signal can be reconstructed. A preliminary characterization of the system is needed to associate the measured duty cycle values to vibration accelerations for each vibration frequency. This way, the vibration acceleration can be directly related to  $D$ .

It is worth mentioning that some approximations were made: when  $V_{gs}$  is equal to the transistor threshold voltage, the resistance between the drain and source is still high. The voltage drop across it is significant, and the diode is still reverse-biased, which implies that the frequency doubler is not in the ON state yet. As  $V_{gs}$  increases, this resistance decreases, and also the voltage potential developed across the resistance decreases. Therefore, a transition region is expected, where  $P_{out}$  features intermediate values between the ON and the OFF states. The threshold voltage  $V_{th}$  that separates the ON from the OFF state of the frequency doubler is defined as the gate–source voltage which minimizes the error caused by the approximation of the measured output power  $P_{out,meas}$  with the theoretical step function  $P_{out,th}$ . Assuming  $P_{out}$  is evaluated for a number  $N$  of gate–source voltage samples from the OFF to the ON state at steps  $\Delta V_{gs}$ , the threshold voltage

$V_{th}$  is defined as follows:

$$V_{th} = \arg \min_{V_{gs}} \text{err}_{P_{out}} \quad (12)$$

where

$$\text{err}_{P_{out}} = \frac{\sum_{n=1}^N |P_{out,dB,meas}(n) - P_{out,dB,th}(n)| \Delta V_{gs}}{N}. \quad (13)$$

As a result, the threshold voltage defined to separate the ON and OFF states of the frequency doubler does not correspond exactly to the threshold voltage of the transistor, and it is actually slightly higher than that.

A significant advantage of the proposed approach is that the whole information on the vibration (i.e., both frequency and acceleration) can be retrieved from the analysis of just the carrier and the first two sidebands of the spectrum of the output signal. However, it is worth mentioning that since the vibration simply causes the amplitude modulation of the backscattered second harmonic, the sensor information can be retrieved also using conventional AM receivers tuned to  $2f_0$ . The proposed method is specifically developed to retrieve frequency and acceleration of sinusoidal vibrations. The estimate of acceleration can be inaccurate if  $v_{gs}$  is distorted.

### III. EXPERIMENTAL RESULTS

A proof-of-concept prototype of the frequency doubler in microstrip technology was designed and manufactured. The adopted substrate was the Rogers RO4350B ( $\epsilon_r = 3.66$ ,  $\tan\delta = 0.0031$ , substrate thickness =  $762 \mu\text{m}$ ). The following off-the-shelf components were used: HSMS2850 from Broadcom was used as the Schottky diode, BSS123 was used as the NMOS transistor, and S129-H5FR-1803YB piezoelectric bending transducer from MIDE was used as the piezoelectric cantilever ( $C_p = 22 \text{ nF}$ ). The doubler was designed with the CAD Advanced Design System using Harmonic Balance simulations for a fundamental frequency  $f_0 = 1.5 \text{ GHz}$  and is optimized to achieve the minimum conversion loss when the inductor  $L_m$  is connected to ground for an available input power in the range  $[-10, -20] \text{ dBm}$ . The SPICE model used for the Schottky diode is provided by the manufacturer and includes package parasitics. No reliable SPICE models of the BSS123 were found, so we used an experimental approach. A photograph of the prototype is shown in Fig. 4.

First, we measured the variation in the power at the RF output port of the frequency doubler  $P_{out}$  with respect to the voltage applied to the gate of the transistor  $V_{gs}$ . We connected the input port of the frequency doubler to an RF signal generator, the output port to a spectrum analyzer, while  $V_{gs}$  was varied with a power supply.

The results, obtained for two different values of the available input power  $P_{in}$ , are shown in Fig. 5. We can identify two different operating regions of the frequency doubler where  $P_{out}$  is constant: an OFF region for  $V_{gs} < 1.4 \text{ V}$  and an ON region for  $V_{gs} > 1.8 \text{ V}$ . For  $P_{in} = -10 \text{ dBm}$ , the conversion loss is equal to 15 dB in the ON region and equal to 31 dB in the OFF region, while for  $P_{in} = -20 \text{ dBm}$ , the conversion loss is equal to 18 dB in the ON region and equal to 31 dB in the OFF region. The threshold voltage of the transistor is equal to 1.4 V.

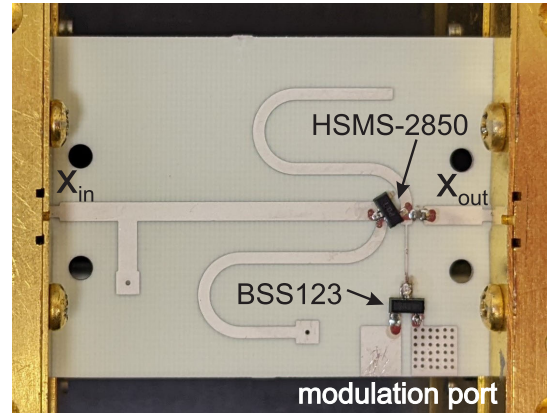


Fig. 4. Photograph of the prototype of the designed frequency doubler for vibration sensing.

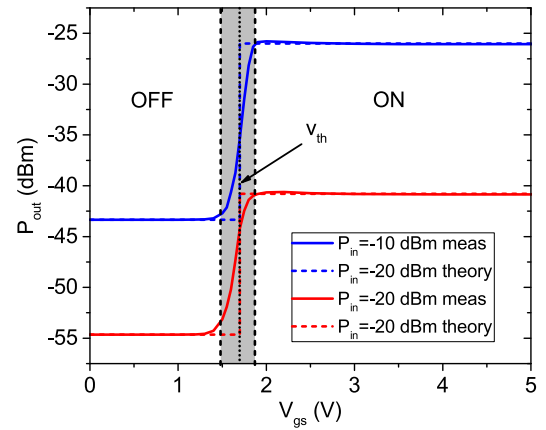


Fig. 5. Measured RF output power versus gate-to-source voltage for different RF input powers.

Between the ON and OFF states, we can see a narrow transition region extending from 1.4 to 1.8 V (in gray in Fig. 5), where the RF output power features intermediate values between the ON and OFF states. As mentioned in Section II, this is due to the fact that at the border of the triode region the resistance measured between the drain and source is still high, and the diode is still reverse-biased. As  $V_{gs}$  increases, the drain–source resistance decreases, and also the voltage potential developed across the resistance decreases, thereby increasing  $P_{out}$ . Since the graph in Fig. 5 is approximated by an ideal step in the developed theory, we define the threshold voltage that separates the ON from the OFF region of the frequency doubler as the gate–source voltage which minimizes the error caused by the step approximation, as described in (13). Therefore,  $V_{th}$  is set to 1.65 V.

We can observe that the RF output power increases by 15 dB from the OFF to the ON state for  $P_{in} = -10 \text{ dBm}$ , and it increases by 12 dB for  $P_{in} = -20 \text{ dBm}$ . The reduced difference between the two states for lower RF input powers is due to the fact that as the RF input power  $P_{in}$  decreases, the dc voltage that is self-generated by the diode  $v_{D,dc}$  is smaller, as shown by (1). Therefore, it has a smaller impact on the frequency doubler performance [34]. Nevertheless, this difference is still considerable for RF input powers as low as

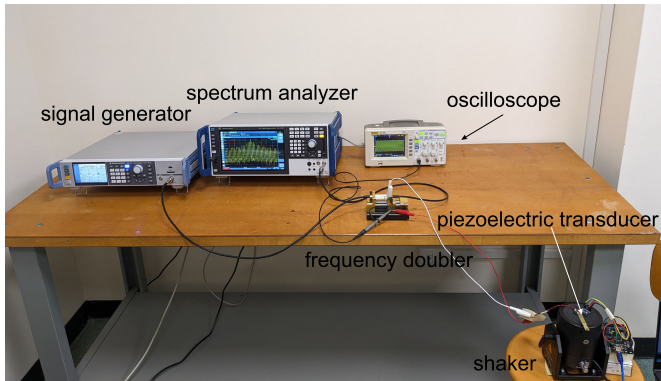


Fig. 6. Proof-of-concept experimental setup for the testing of the frequency-doubler-based sensor using a piezoelectric cantilever.

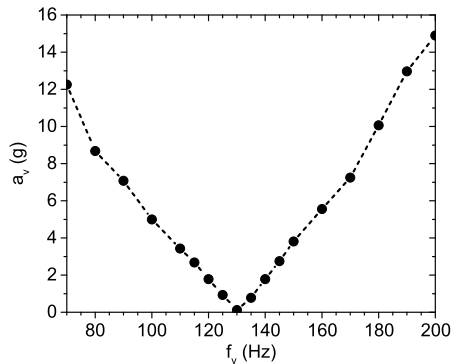


Fig. 7. Cantilever characterization: acceleration corresponding to  $V_{gs} = 2$  V for different vibration frequencies.

−20 dBm, which guarantees a high transponder sensitivity for wireless applications. Indeed, assuming  $f_0 = 2.45$  GHz (ISM band), in case the power transmitted by the reader is 36 dBm, this sensitivity corresponds to a read range of 9 m.

The setup adopted to test the frequency doubler with a piezoelectric cantilever is shown in Fig. 6. The transistor is connected to the piezoelectric cantilever using two wires. One side of the cantilever is fixed to a shaker (model DS-PM-20 from DEWESoft) with a screw. A commercial accelerometer (model ADXL354 from Analog Devices) is also placed on the shaker next to the cantilever to measure the acceleration of the applied vibration. An oscilloscope voltage probe is connected in parallel with the cantilever to check the voltage waveforms generated by the cantilever in real-time.

First, we characterized the adopted piezoelectric cantilever. The natural oscillation frequency of the cantilever is estimated as the vibration frequency where the lowest acceleration is needed to achieve a certain generated voltage amplitude. The vibration frequency is varied in the range of 70–200 Hz at steps of 5 Hz. The results are shown in Fig. 7 for  $V_{gs} = 2$  V (i.e., just above threshold). Based on the results, the estimated natural oscillation frequency is 130 Hz. In this study, we aimed at achieving the highest vibration sensitivity. Therefore, we used the transducer around its resonance.

We analyzed the relationship between the RF signal at the output port of the frequency doubler and the gate-to-source voltage generated by the piezoelectric transducer. The

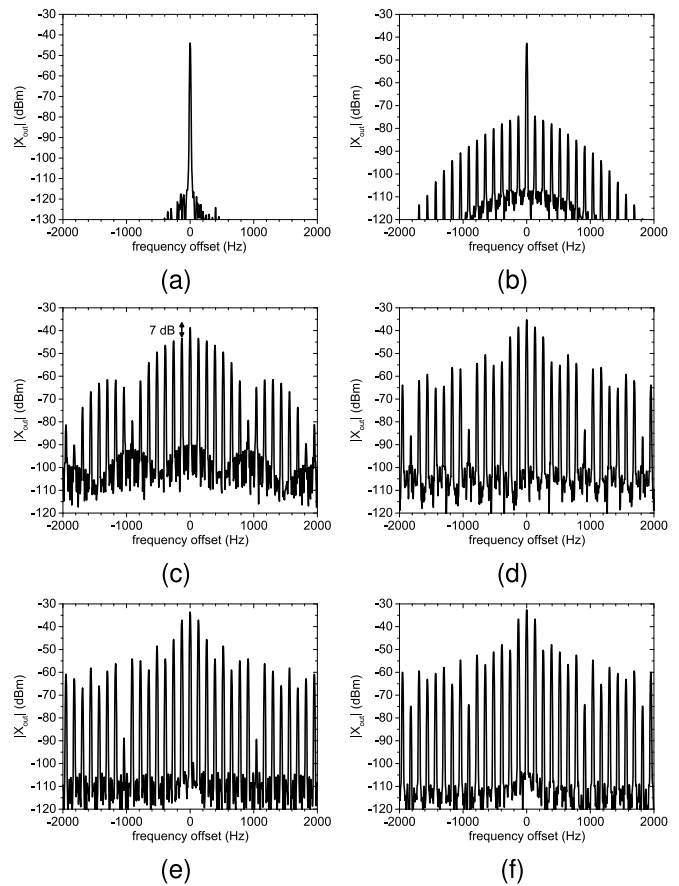


Fig. 8. Measured output spectrum for  $f_v = 130$  Hz and  $P_{in} = -10$  dBm. (a)  $V_{gs} = 0$  V (corresponding to no vibration), (b)  $V_{gs} = 1.6$  V, (c)  $V_{gs} = 1.9$  V, (d)  $V_{gs} = 2.45$  V, (e)  $V_{gs} = 4.3$  V, and (f)  $V_{gs} = 8$  V.

vibration frequency of the shaker is set to 130 Hz. The transmitted RF input power is set to −10 dBm, and the vibration acceleration is varied so that a variable  $V_{gs}$  is generated by the piezoelectric cantilever. Sample spectra acquired at the RF output port for different accelerations are shown in Fig. 8. The acquisition band of the spectrum analyzer is set to 4 kHz, the resolution band is 10 Hz, and the number of points is 801. These parameters are set to achieve a satisfactory resolution, low noise floor, while keeping fast sweep time.

A simple MATLAB code is developed to retrieve the vibration information from the spectrum. The code first identifies local maxima in the acquired spectrum using the function `findpeaks`. To avoid false peaks due to noise, all the peaks which are below −100 dBm are disregarded by defining “MinPeakHeight,” while the minimum distance between two peaks is set to 60 Hz by defining “MinPeakDistance.” The “MinPeakHeight” value is determined heuristically based on the measured noise floor of the spectrum analyzer, while the “MinPeakDistance” value is determined based on the operating frequency band of the used piezoelectric transducer. Machine-learning-based estimators can also be used to reduce the risk of errors in noisy signal acquisitions.

If there are no vibrations, no sidebands are present in the spectrum, as shown in Fig. 8(a). If the vibrations are too

weak to generate a voltage above threshold, there might be sidebands in the spectrum, but they are well below the carrier (an example is shown in Fig. 8(b) where  $V_{gs} = 1.6$  V and the first sideband is about 20 dB below the carrier). Finally, if the vibrations are strong enough to generate  $V_{gs}$  greater than  $V_{th}$ , sidebands appear in the spectrum according to the proposed theory. In Fig. 8(c), the first sidebands  $A_1$  are about 7 dB below the carrier and this case can be considered borderline, since  $V_{gs}$  is slightly higher than  $V_{th}$ . Therefore, for a given vibration frequency, the ratio of the carrier component to the first sideband is computed and the acceleration corresponding to  $A_{01} = A_0/A_1 = -7$  dB is set as the lowest acceleration that can be measured by the transponder. As the vibration acceleration increases, the voltage amplitude generated by the piezoelectric cantilever increases as well (the measured values are reported in the caption of Fig. 8), which causes a larger  $D$ . Consequently, the amplitude of the carrier increases, since it depends linearly on  $D$ , as shown in (9). In addition, the relative amplitude of the sidebands varies.

The vibration frequency is calculated as the average frequency separation between the first sidebands and the carrier tone. Then, the duty cycle  $D$  of the waveform is determined from  $A_{21}$  as described in (11). To verify the consistency of the process, the signal envelope is also recovered. The Fourier coefficients are calculated as described in (7). Finally, the time-domain waveform is retrieved using (5).

The duty cycle of the modulating signal, retrieved from the acquired spectrum using the parameter  $A_{21}$  (“measurements”), is then compared with the duty cycle calculated using (3) (“theory”) based on  $V_{gs}$  amplitude directly measured with the oscilloscope, as shown in Fig. 9(a). We observe an excellent agreement between the measurements and theory as long as  $V_{gs}$  is larger than  $V_{th}$ .

The retrieved time-domain modulation signals are shown in Fig. 9(b) and (c) for two different  $V_{gs}$ . Square waves with different duty cycles can be clearly identified in both the cases. The retrieved RF signal envelope is confronted with the respective voltage signal  $v_{gs}$  generated by the piezoelectric cantilever. We can observe the close correspondence between the fraction of the vibration period in which the voltage signal is above  $V_{th}$  and the duty cycle of the signal envelope. The modulation frequency measured with the above-described method is equal to 130 Hz for all the spectra, in perfect agreement with the nominal vibration frequency displayed by the shaker.

The RF output signal is also acquired for different values of the RF input power, for a fixed vibration frequency ( $f_v = 130$  Hz). The results are illustrated in Fig. 10. We can observe a good agreement in the duty cycle estimation in all the cases in Fig. 10(a), while the estimated modulation frequency is the same in all the cases (equal to 130 Hz), as shown by the reconstructed time-domain modulation signals in Fig. 10(b) for  $V_{gs} = 7.4$  V.

Then, we analyzed the relationship between the duty cycle of the RF output signal of the frequency doubler and the vibration acceleration. All the piezoelectric transducers are characterized by a frequency response which depends on their shape, weight, and elasticity [44]. Therefore, a calibration step

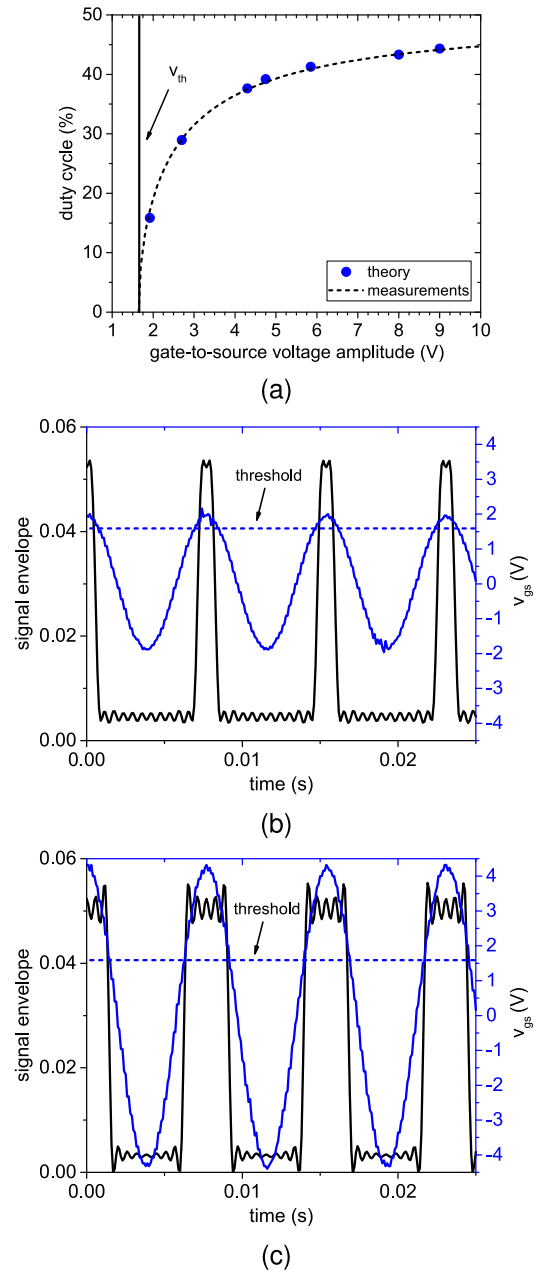
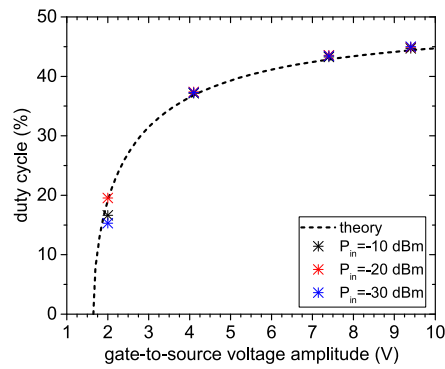
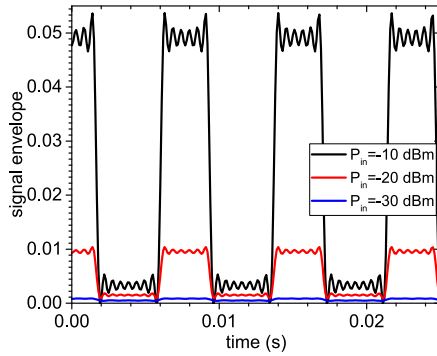


Fig. 9. RF output signals at  $f_v = 130$  Hz for different vibration accelerations. (a) Measured duty cycle, (b) retrieved modulation signal and measured gate-to-source voltage for  $V_{gs} = 1.9$  V, and (c) for  $V_{gs} = 4.3$  V.

is needed to determine the relationship between the measured duty cycle and the vibration acceleration for each vibration frequency. To identify the operating frequency band of the vibration sensor, we set the vibration acceleration to 2 g, and we varied the vibration frequency of the shaker. The measured duty cycles versus vibration frequency are shown in Fig. 11. The tag is able to detect vibrations in the frequency range from 115 to 142 Hz for a vibration acceleration of 2 g. The highest duty cycle is measured at the resonance frequency of the cantilever, namely, at 130 Hz, where the piezoelectric cantilever generates the largest  $V_{gs}$ . It is worth mentioning that the operating frequency band of the proposed vibration sensor depends on the adopted piezoelectric cantilever, since



(a)



(b)

Fig. 10. RF output signals at  $f_v = 130$  Hz for different RF input powers. (a) Measured  $D$  and (b) retrieved envelopes of the measured modulated signals for  $V_{gs} = 7.4$  V.

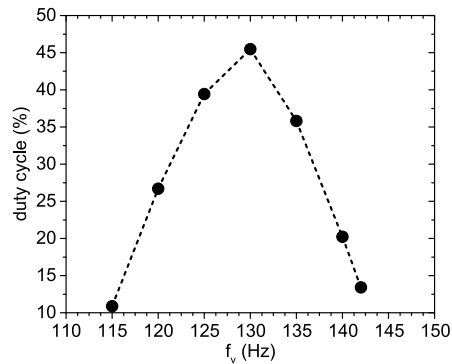
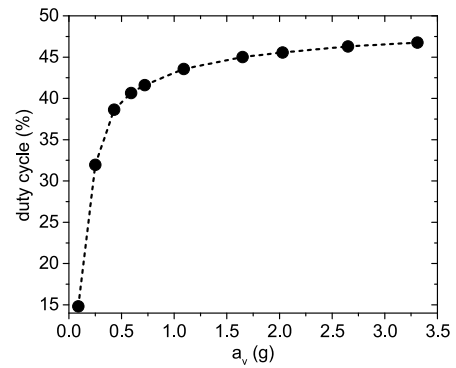


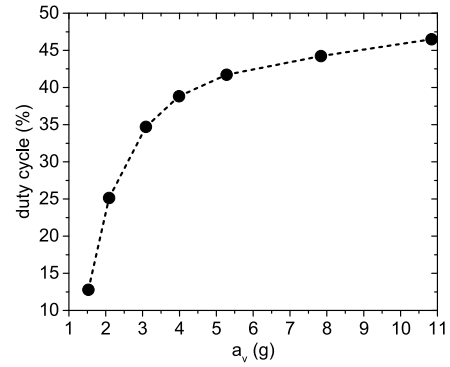
Fig. 11. Measured duty cycle versus vibration frequency (acceleration 2 g).

the proposed circuit can sustain modulation frequencies up to 10 MHz, as already demonstrated in [34]. Therefore, transducers with different natural oscillation frequencies can be used to make the sensor sensitive to different vibration frequencies.

In Fig. 12(a) and (b), the duty cycle of the RF output signal versus the vibration acceleration are shown for two vibration frequencies. In the first case, the vibration frequency of the shaker is set to 130 Hz, and the vibration acceleration is varied. As shown in Fig. 12(a), the measured duty cycle parameter varies from about 14.8% to 46.8% for a vibration acceleration ranging from 0.09 to 3.3 g. The estimated vibration frequency with the developed MATLAB code is 130 Hz in all the cases. The same experiment is repeated for  $f_v = 120$  Hz,



(a)



(b)

Fig. 12. Measured duty cycle versus vibration acceleration for (a)  $f_v = 130$  and (b)  $f_v = 120$  Hz.

and the results are shown in Fig. 12(b). The measured duty cycle parameter varies from about 12.8% to 46.5% for a vibration acceleration ranging from 1.5 to 10 g. In both the cases, the relationship  $D-a_v$  is nonlinear, and  $D$  increases monotonically as  $a_v$  increases. The largest vibration frequency that can be sensed by the transducer is limited by the saturation of  $D(a_v)$  function. The nonlinear  $D-a_v$  relationship makes this transponder a good candidate for shock sensing. Acceleration sensitivity increases significantly close to the natural oscillation frequency of the piezoelectric cantilever, since, for a given acceleration, the cantilever deflection is larger than at any other oscillation frequency, which means that a larger voltage signal is generated. The sensitivity progressively degrades moving away from the resonance, although the acceleration sensing range significantly increases. The transducer should be used at resonance to detect weak vibrations and off resonance to monitor larger accelerations.

These results demonstrate the high sensitivity of the proposed solution and its suitability to perform condition monitoring for electrical machines and industrial equipment [12]. The threshold voltage of the MOSFET impacts on the sensitivity of the sensor, since, for a given piezoelectric transducer, a lower threshold voltage allows the sensor to measure weaker vibrations. If the threshold voltage of the FET is unknown, a calibration step is needed to accurately measure vibration accelerations.

Finally, a complete wireless fully passive sensor—with wireless interrogation capability—was manufactured and tested. The described modulation-capable frequency doubler



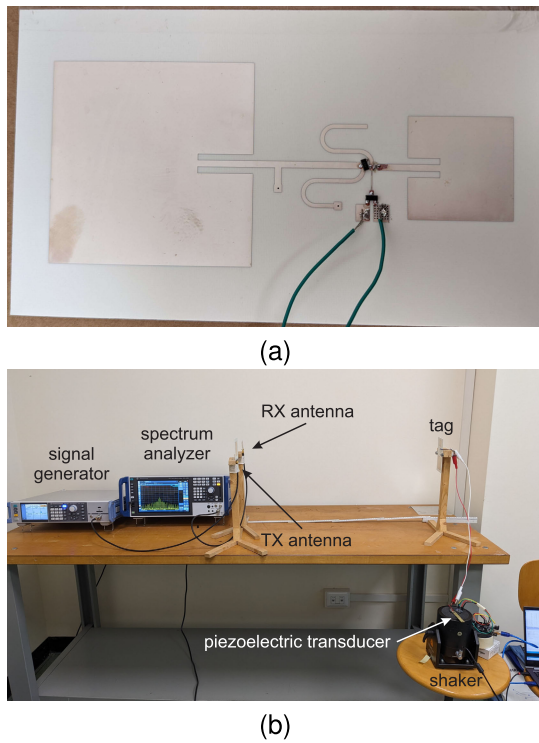


Fig. 13. Test of the wireless vibration sensor: (a) photograph of the tag and (b) experimental setup.

is connected to an antenna working at  $f_0 = 1.5$  GHz at the input port, and to an antenna working at  $2f_0$  at the output port. A photograph of the manufactured proof-of-concept prototype is shown in Fig. 13(a). Two square patch antennas are used as the radiating elements. The two antennas are matched to  $50 \Omega$ , and the whole prototype is manufactured on Rogers 4350B. The overall area of the PCB is  $13.5 \times 8$  cm<sup>2</sup>.

The adopted test setup is shown in Fig. 13(b). The signal generator is connected to a patch antenna with its first resonance at  $f_0$  and the spectrum analyzer is connected to a patch antenna with its first resonance at  $2f_0$  (the tag and the reader antennas are the same in the adopted proof-of-concept setup). The power transmitted at  $f_0$  is 15 dBm, and the same settings as in the previous experiment are used for the spectrum analyzer (namely, acquisition band 4 kHz, resolution band 10 Hz). The threshold power to avoid spurious peaks in the MATLAB code is lowered to  $-120$  dBm due to the lower noise floor experienced in the wireless experiment due to the weaker received signal.

The vibration frequency of the shaker is set to 130 Hz, and, for each vibration acceleration, the tag-to-reader distance is varied from 50 to 130 cm. The tag-to-reader distance is chosen to have an RF power  $P_{in}$  at the input of the doubler in the design range  $[-10, -20]$  dBm. The coverage can be further improved by increasing the transmitted power and using an LNA at the receiver input. The measured spectra are illustrated in Fig. 14 for  $a_v = 1$  g. As expected, the received power varies depending on the tag-to-reader distance, on the multipath, and so forth. However, the carrier and the first two sidebands, which are needed to recover the sensor information, are still significantly above the noise floor, and the magnitude ratio of

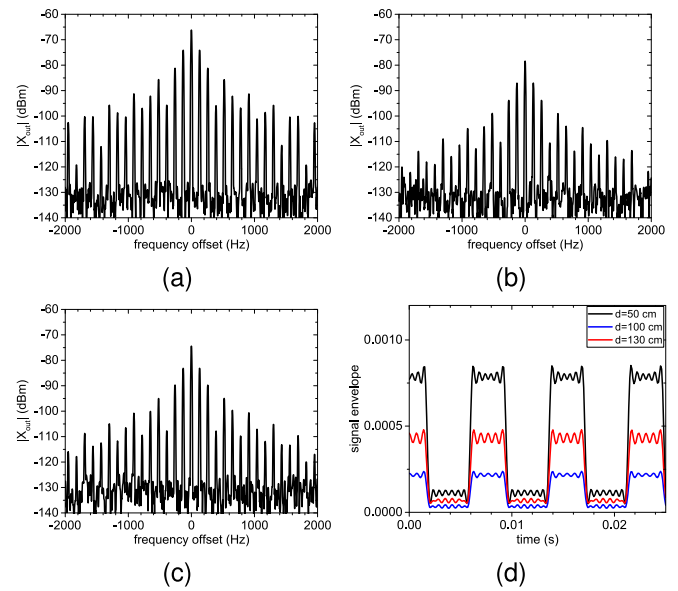


Fig. 14. Received signals versus tag-to-reader distance for  $f_v = 130$  Hz and  $a_v = 1$  g. (a) Spectrum for  $d = 50$  cm, (b) spectrum for  $d = 100$  cm, (c) spectrum for  $d = 130$  cm, and (d) reconstructed time-domain signal envelope versus distance.

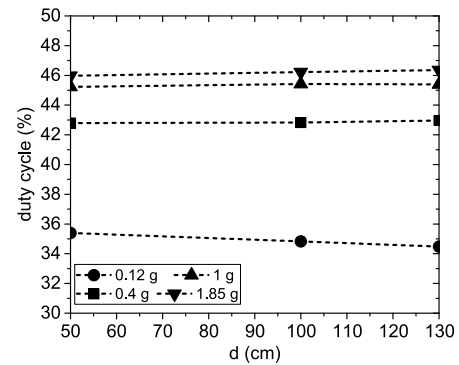


Fig. 15. Measured duty cycle at  $f_v = 130$  Hz for different tag-to-reader distances and vibration accelerations.

the first and second sidebands is the same, except for the noise of the acquisition instrument.

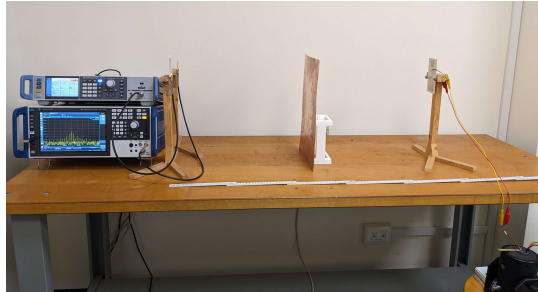
The estimated vibration frequency is equal to 130 Hz for all the measurements. Fig. 14(d) shows the reconstructed envelope signals. The square waves are clearly recognizable for all the three distances.

The estimated duty cycle versus distance at  $f_v = 130$  Hz is shown in Fig. 15 for different vibration accelerations. We observe that for a given acceleration, the measured duty cycle is almost the same regardless of the tag-to-reader distance. Small fluctuations are due to noise. The noise impact can be mitigated by averaging repeated measurements.

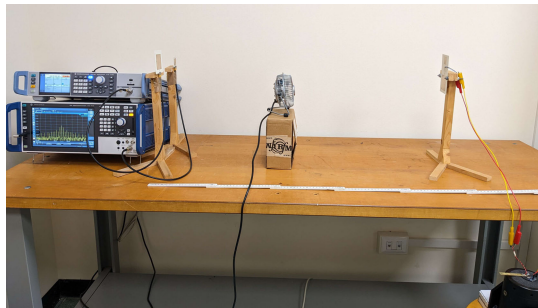
Finally, we tested the capability of the tag to operate in highly cluttered environments and in the presence of dynamic multipath, to test its suitability for industrial environments. Since vibrations induce the amplitude modulation of the backscattered second harmonic, multipath fading can cause the attenuation of single spectral components, thus causing errors in the vibration acceleration estimation. To mitigate this, the estimation of the duty cycle is made by averaging

TABLE I  
COMPARISON TO THE STATE OF THE ART

Ref.	passive	transponder type	transducer type	measured parameters	vibr. frequency (Hz)	vibr. acceleration (g)
this work	yes	harmonic tag	piezoelectric cantilever	acceleration and frequency	115 – 143	above 2
[25]	no	antenna backscattering	strain gauge	strain	120	1.5 – 11
[26]	no	antenna backscattering	microstrip patch	strain	130	0.09 – 3.5
[28]	yes	LC resonator	capacitor	acceleration	11.9	/
[27]	yes	SAW device	piezoelectric cantilever	acceleration and frequency	1-8	/
[29]	yes	harmonic tag	piezoelectric cantilever	acceleration and frequency	120	0.4 – 5.1
[29]	yes	harmonic tag	piezoelectric cantilever	acceleration and frequency	20	0-2.5
[29]	yes	harmonic tag	mechanical switch	frequency	5-85	9-40
					0.1	0.01
					30	2.3



(a)



(b)

Fig. 16. Experimental setup to test the tag in an environment with clutter and multipath: (a) metal panel between the tag and the reader and (b) rotating fan between the tag and the reader.

the ratio obtained with the right-hand sidebands with the one obtained with the left-hand sidebands. Concerning dynamic multipath, as long as it is slower than the vibration frequency, it does not significantly affect the communication. Fig. 16 shows the experimental setups. First, we put a metal panel ( $30 \times 50 \text{ cm}^2$ ) between the reader antennas and the tag (absence of line of sight and strong  $f_0$  reflections), and then we replace the metal panel with a rotating fan. The tag-to-reader distance is 100 cm. The acquired spectrum and measured duty cycle are shown in Fig. 17(a) and (b), respectively. Vibration frequency and acceleration are kept constant in all the experiments. Although attenuated, the acquired spectrum preserves its shape in presence of both types of obstacles, thus confirming the capability of the sensor to operate in harsh environments.

In the present work, we tested the capability of the system to recover the vibration information from a single wireless passive tag. In case multiple transponders are used, they can be separated in the frequency domain, by designing each tag

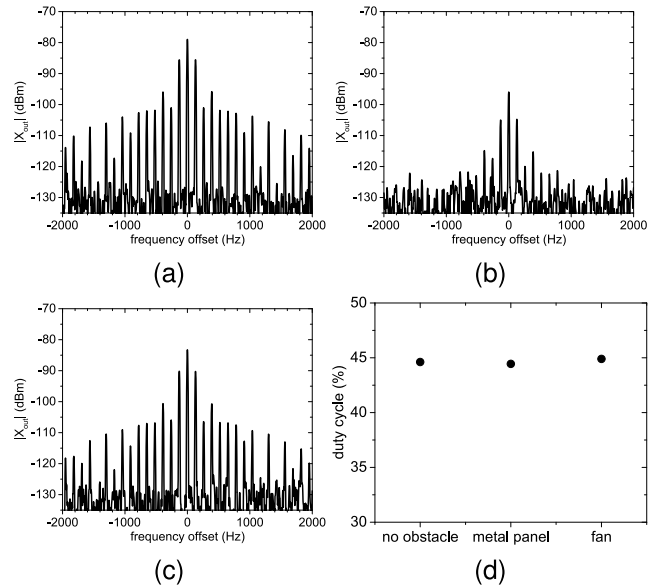


Fig. 17. Experimental results of the wireless vibration sensing system in the presence of obstacles. (a) Acquired spectrum without any obstacle, (b) acquired spectrum with a metal panel between the tag and the reader, (c) acquired spectrum with a rotating fan between the tag and the reader, and (d) measured duty cycle.

for different RF frequencies and by making the transmitter sweep over frequency.

The presented wireless passive vibration sensor is then compared with other state-of-the-art wireless vibration sensors based on backscatter radios in Table I. The proposed solution features an excellent compromise among low-cost, sensitivity, and versatility, while representing a fully passive wireless solution, and the obtained results open the door to a novel acquisition method for vibration sensing in industrial environments.

#### IV. CONCLUSION

A novel vibration sensing approach has been presented, based on a passive harmonic transponder and a piezoelectric cantilever. It has been shown that the piezoelectric transducer is able to modulate the amplitude of the harmonic backscattered signal by controlling the gate of an NMOS switch. The vibration frequency is encoded in the frequency difference between the spectral lines of the received signal, while the acceleration is encoded in the duty cycle of the amplitude modulation of the backscattered signal. A method to derive the

sensor information using only the spectrum of the harmonic backscattered signal (scalar noncoherent measurement) has been described in an implementation that does not depend on the tag-to-reader distance. A prototype has been manufactured and tested, and the concept has been proven successfully. This work opens the door to a new class of harmonic wireless passive sensors suitable for industrial applications.

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